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Fully Automated Calculations in the complex MSSM

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We review recent progress towards automated higher-order calculations in the MSSM with complex parameters (cMSSM). The consistent renormalization of all relevant sectors of the cMSSM and the inclusion into the FeynArts/FormCalc framework has recently been completed. Some example calculations applying this framework are briefly discussed. These include two-loop corrections to cMSSM Higgs boson masses as well as partial decay widths of electroweak supersymmetric particles decaying into a Higgs boson and another supersymmetric particle.

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1. Introduction

Two of the most important goals of the experiments at the Large Hadron Collider (LHC) are to identify the origin of electroweak symmetry breaking (EWSB), and to search for physics effects beyond the Standard Model (SM). The spectacular discovery of a Higgs-like particle with a mass around ~ 125.6 GeV, which was announced by ATLAS and CMS [1, 2], marks a milestone of an effort that has been ongoing for almost half a century and opens a new era of particle physics. Within the experimental uncertainties the properties of the newly discovered particle are in agreement with the predictions of the SM Higgs boson [3, 4]. However, the uncertainties still leave room for contributions from non-SM degrees of freedom, see, e.g., Refs. [5–7] for a recent combination and reviews. The prime task now is to study the properties of the discovered new particle in detail and to investigate whether there are significant deviations from the SM predictions, which would point towards physics beyond the SM.

The extent to which the results of the Higgs searches at the LHC can discriminate between the SM and possible alternatives depends both on the experimental precision with which the properties of a possible signal can be determined and on the detailed nature of the mechanism for EWSB that is actually realized in nature. One of the leading candidates for physics beyond the SM (BSM) is supersymmetry (SUSY), which doubles the particle degrees of freedom by predicting two scalar partners for all SM fermions, as well as fermionic partners to all bosons. The most widely studied SUSY framework is the Minimal Supersymmetric Standard Model (MSSM) [8], which keeps the number of new fields and couplings to a minimum. The MSSM Higgs sector contains two Higgs doublets, which at the tree-level leads to a physical spectrum consisting of two $\mathcal{C}\mathcal{P}$ -even, h, H , one $\mathcal{C}\mathcal{P}$ -odd, A , and two charged Higgs bosons, H^\pm .

In order to investigate the impact of the Higgs search results at the LHC on possible scenarios of new physics, precise theoretical predictions both within the SM and possible alternatives of it are needed. In particular, if small deviations from the SM predictions are probed it is crucial to treat the considered model of new physics at the same level of precision to enable an accurate analysis and comparison. In the MSSM Higgs sector higher-order contributions are known to give numerically large effects (see, e.g., Refs. [9, 10]). For many observables it is therefore necessary to include corrections beyond leading order in the perturbative expansion to obtain reliable results. The calculation of loop diagrams, often involving a large number of fields, is a tedious and error-prone task if done by hand. This is true in particular for BSM theories where the number of fields is significantly increased. For one-loop calculations, as will be the focus in the following, computer methods with a high degree of automatization have been devised to simplify the work. However, most of the available tools so far have focused on calculations either in the SM or the MSSM with external SM particles.

Here we review renormalization of the MSSM including complex parameters (*c*MSSM) and the corresponding implementation as a model file [11] into the `FeynArts` [12, 13]/`FormCalc` [14] framework. This implementation allows for automated calculation of processes with external SUSY particles. We also briefly discuss the application of the new `FeynArts` model file to the calculation of two-loop corrections to *c*MSSM Higgs boson masses [15, 16], and to the evaluation of partial decay widths of SUSY electroweak (EW) particles [17–19].

2. Renormalization of the cMSSM

The tree-level Feynman rules of the MSSM are by now well under control, where the cMSSM had been included into the `FeynArts` package [13]. Concerning the renormalization, however, most calculations in the past chose a prescription that was tailored to one specific calculation or even one specific part of the (c)MSSM parameter space. Since the values of the SUSY parameters realized in nature are unknown, at the current state scans over large parts of the cMSSM parameter space are necessary. Furthermore, many processes have to be evaluated simultaneously. Both requirements make a *complete* renormalization of the cMSSM necessary that is valid over the *full* (or at least “large parts”) of the cMSSM parameter space. Only with such a renormalization at hand fully automated calculations in the cMSSM will be possible. Evidently, calculations at n -loop require an n -loop renormalization, where we will focus on the one-loop case.

The program of the renormalization of all (physical) sectors of the cMSSM has recently been completed [11, 17–24] (based on earlier work in the MSSM [25–28]; for alternative approaches see Refs. [29–31]) and included as a model file `MSSMCT.mod` [11] into the `FeynArts` package.

In the development of the renormalization particular emphasis was put on the requirement that the one-loop corrections stay “small” over the full allowed parameter range. The renormalization includes the scalar fermion sector, the remaining colored sector, the chargino/neutralino sector and the Higgs sector (as well as the SM part of the MSSM). In principle this is sufficient to evaluate all currently relevant processes at the one-loop level. Extensive checks have been performed to ensure “stability” of the higher-order corrections over large(st) parts of the cMSSM parameter space. These tests include scalar top and bottom decays [21, 22], scalar tau decays [24], gluino decays [23] as well as non-hadronic chargino [17] and neutralino decays [18, 19]. These evaluations are complete at the one-loop level, including hard and soft QED and QCD radiation.

To give an idea about the new model file we briefly review the various masses, coupling constants etc. as well as the respective counterterms implemented into `MSSMCT.mod` (all details can be found in Ref. [11]). In Tab. 1 we list the particle content of the model file, where the respective index ranges are given in Tab. 2. The symbols for the masses of the particles, for couplings and mixing angles are shown in Tab. 3. Within the Higgs boson sector the tree-level masses are taken distinct from the higher-order corrected masses (which can be obtained via the automatic link to `FeynHiggs` [20, 32–35]). $\tan\beta$ denotes the ratio of the two vacuum expectation values, and α is the angle that diagonalizes the (tree-level) \mathcal{CP} -even Higgs sector. Within the cMSSM all three neutral Higgs bosons can mix to give rise to three \mathcal{CP} -mixed states, h_i ($i = 1, 2, 3$). When composing a vertex Γ_{h_i} from the corresponding tree-level amplitudes Γ_h , Γ_H , and Γ_A , a set of finite Z -factors is needed to ensure correct on-shell properties of the external Higgs boson h_i [20],

$$\Gamma_{h_i} = \hat{Z}_{i1}\Gamma_h + \hat{Z}_{i2}\Gamma_H + \hat{Z}_{i3}\Gamma_A + \dots, \quad (2.1)$$

where the ellipsis represents contributions from the mixing with the Goldstone and Z boson. The Z -factor matrix $\hat{Z}_{ij} \equiv Z_{\text{Higgs}}[i, j]$ is not in general unitary. Its lower 3×3 part is computed by `FeynHiggs` and application at the amplitude level automatically takes any absorptive contribution into account. Technically this is most easily accomplished using the `FeynArts` add-on model file `HMix.mod` [36] which mixes $h = S[1]$, $H = S[2]$, and $A = S[3]$ into two variants of the

leptons	$f = f^\dagger$	field	mass	sleptons	$f = f^\dagger$	field	mass
ν_g		$F[1, \{g\}]$	0	$\tilde{\nu}_g$		$S[11, \{g\}]$	MSf
ℓ_g		$F[2, \{g\}]$	MLE	$\tilde{\ell}_g^s$		$S[12, \{s, g\}]$	MSf

quarks				squarks			
u_g		$F[3, \{g, o\}]$	MQU	\tilde{u}_g^s		$S[13, \{s, g, o\}]$	MSf
d_g		$F[4, \{g, o\}]$	MQD	\tilde{d}_g^s		$S[14, \{s, g, o\}]$	MSf

gauge bosons				neutralinos, charginos			
γ	yes	$V[1]$	0	$\tilde{\chi}_n^0$	yes	$F[11, \{n\}]$	MNeu
Z	yes	$V[2]$	MZ	$\tilde{\chi}_c^-$		$F[12, \{c\}]$	MCha
W^-		$V[3]$	MW				

Higgs/Goldstone bosons				ghosts			
h	yes	$S[1]$	Mh0	u_γ		$U[1]$	0
H	yes	$S[2]$	MHH	u_Z		$U[2]$	MZ
A	yes	$S[3]$	MA0	u_+		$U[3]$	MW
G	yes	$S[4]$	MZ	u_-		$U[4]$	MW
H^-		$S[5]$	MHp	u_g		$U[5, \{u\}]$	0
G^-		$S[6]$	MW				

gluon				gluino			
g	yes	$V[5, \{u\}]$	0	\tilde{g}	yes	$F[15, \{u\}]$	MGl

Table 1: The particle content of MSSMCT.mod.

$g = \text{Index}[\text{Generation}] = 1 \dots 3,$
 $o = \text{Index}[\text{Colour}] = 1 \dots 3,$
 $u = \text{Index}[\text{Gluon}] = 1 \dots 8,$
 $s = \text{Index}[\text{Sfermion}] = 1 \dots 2,$
 $n = \text{Index}[\text{Neutralino}] = 1 \dots 4,$
 $c = \text{Index}[\text{Chargino}] = 1 \dots 2.$

(S)fermions are indexed by

$$t = \begin{cases} 1 & \text{(s)neutrinos,} \\ 2 & \text{charged (s)leptons,} \\ 3 & \text{up-type (s)quarks,} \\ 4 & \text{down-type (s)quarks.} \end{cases}$$

Table 2: Index labels and ranges used in MSSMCT.mod.

Mh0, MHH, MA0, MHp Mh0tree, MHHtree, MA0tree, MHptree TB, CB, SB, C2B, S2B CA, SA, C2A, S2A CAB, SAB, CBA, SBA MUE	Higgs masses M_h, M_H, M_A, M_{H^\pm} tree-level Higgs masses $\tan\beta, \cos\beta, \sin\beta, \cos 2\beta, \sin 2\beta$ $\cos\alpha, \sin\alpha, \cos 2\alpha, \sin 2\alpha$ (tree-level α) $\cos(\alpha + \beta), \sin(\alpha + \beta), \cos(\beta - \alpha), \sin(\beta - \alpha)$ Higgs-doublet mixing parameter μ
MGl SqrtEgl MNeu[n] ZNeu[n, n'] MCha[c] UCha[c, c'], VCha[c, c']	gluino mass $m_{\tilde{g}}$ root of the gluino phase, $e^{i\phi_{\tilde{g}}/2}$ neutralino masses $m_{\tilde{\chi}_n^0}$ neutralino mixing matrix $\mathbf{N}_{nn'}$ chargino masses $m_{\tilde{\chi}_c^\pm}$ chargino mixing matrices $\mathbf{U}_{cc'}, \mathbf{V}_{cc'}$
MSf[s, t, g] USf[t, g][s, s'] Af[t, g, g']	sfermion masses $m_{\tilde{f}_{i,sg}}$ sfermion mixing matrix $U_{ss'}^{\tilde{f}_{ig}}$ soft-breaking trilinear A-parameters $(\mathbf{A}_{f_i})_{gg'}$
MW, MZ Mf[t, g]	gauge-boson masses M_W, M_Z fermion masses $m_{f_{ig}}$
CW, SW EL GS	$c_w \equiv \cos\theta_w = M_W/M_Z, s_w \equiv \sin\theta_w$ electromagnetic coupling constant e strong coupling constant g_s

Table 3: Symbols representing the SM and MSSM parameters in MSSMCT.mod.

loop-corrected states h_i ,

$$S[0, \{i\}] = \sum_{j=1}^3 \text{UHiggs}[i, j] S[j], \quad \text{with unitary UHiggs (no absorptive contrib.), for use on internal lines,} \quad (2.2a)$$

$$S[10, \{i\}] = \sum_{j=1}^3 \text{ZHiggs}[i, j] S[j], \quad \text{inserted only on external lines.} \quad (2.2b)$$

More details can be found in Refs. [11, 20]. The renormalization constants (RCs) are listed in Tab. 4. For the Higgs boson sector, besides the mass and field renormalization constants also the tadpole counterterms are listed, which correspond to the terms linear in the Higgs fields in the Higgs potential. Several field renormalization constants are given in a barred and unbarred version, differentiating between incoming/outgoing (anti)particles, see Refs. [11, 17, 18, 22] for more details. Several electroweak SM RCs have been omitted (such as dMW1 etc.), since they are defined identical to the pure SM case. The FeynArts/FormCalc framework provides a default implementation of the determination of all RCs, where again the details can be found in Refs. [11, 17–24]. It should be stressed again that this default implementation is based on the requirement of “stability” of the higher-order corrections over large(st) parts of the cMSSM parameter space.

Higgs-boson Sector	
$dZ[\text{bar}]Higgs1[h, h']$	Higgs field RCs
$dMHiggs1[h, h']$	Higgs mass RCs
$dTh01, dTHH1, dTA01$	Higgs tadpole RCs
$dZH1, dZH2, dTB1, dSB1, dCB1$	RCs related to β
Gauge-boson Sector	
$dMZsq1, dMWSq1$	gauge-boson mass RCs
$dZAA1, dZAZ1, dZZA1, dZZZ1, dZ[\text{bar}]W1$	gauge-boson field RCs
$dSW1, dZe1$	coupling-constant RCs
Chargino/Neutralino Sector	
$dMCha1[c, c']$	chargino mass RCs
$dMNeu1[n, n']$	neutralino mass RCs
$dMino11, dMino21, dMUE1$	RCs for M_1, M_2, μ
$dZ[\text{bar}]f_R^L[12, c, c']$	chargino field RCs
$dZ[\text{bar}]f_R^L[11, n, n']$	neutralino field RCs
Fermion Sector	
$dMf1[t, g]$	fermion mass RCs
$dZ[\text{bar}]f_R^L[t, g, g']$	fermion field RCs
$dCKM1[g, g']$	CKM-matrix RCs
Squark Sector	
$dMSfsq1[s, s', 3 4, g]$	squark mass RCs
$dAf1[3 4, g, g]$	trilinear squark coupling RCs
$dZ[\text{bar}]Sf_R^L[s, s', 3 4, g]$	squark field RCs
Slepton Sector	
$dMSfsq1[s, s', 1 2, g]$	slepton mass RCs
$dAf1[2, g, g]$	trilinear slepton coupling RCs
$dZ[\text{bar}]Sf_R^L[s, s', 1 2, g]$	slepton field RCs
Gluino Sector	
$dMG11$	gluino mass RC
$dZ[\text{bar}]G1_R^L$	gluino field RCs
Gluon Sector	
$dZgs1$	strong-coupling-constant RC
$dZGG1$	gluon field RCs

Table 4: RCs used in `MSSMCT.mod`, where $a|b$ means ‘ a or b ’ and $dZ[\text{bar}]$ stands for both dZ and $dZ\text{bar}$ (see text).

3. Example applications

In this section we briefly review some of the yet existing applications of `MSSMCT.mod`.

3.1 Two-loop corrections to Higgs boson masses

Two-loop corrections to cMSSM Higgs boson masses, obtained in the Feynman-diagrammatic approach, require the calculation of two-loop Higgs-boson self-energies, which in turn require a renormalization of the Higgs boson sector at the two-loop level (see Refs. [15, 16, 20, 33, 37–39]). Furthermore, a sub-loop renormalization of the corresponding one-loop diagrams is necessary. Consequently, a full two-loop calculation of the Higgs-boson self-energies requires a full renormalization of the cMSSM at the one-loop level. Recently, two new two-loop calculations were presented. One consists of the $\mathcal{O}(\alpha_t^2)$ contributions involving complex phases [16], the other one of the momentum dependent two-loop part of the $\mathcal{O}(\alpha_t \alpha_s)$ corrections for real parameters [15]. Both types of corrections can yield contributions to M_h larger than the current experimental uncertainty. Corresponding one-loop diagrams with sub-loop renormalization are depicted in Fig. 1 (taken from Ref. [15]).

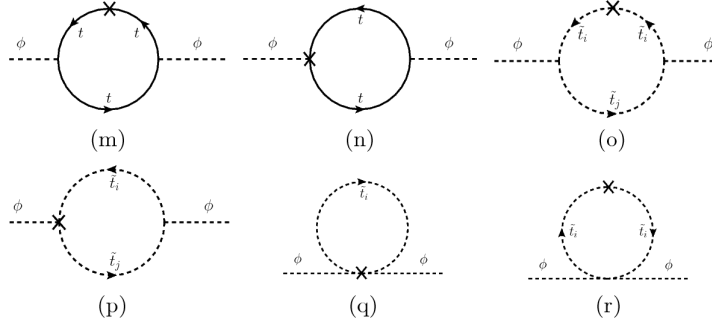


Figure 1: Generic one-loop contributions to cMSSM Higgs-boson self-energies with sub-loop renormalization; $\phi = h, H, A$; t denotes the top quark; $\tilde{t}_{i,j}$ the scalar tops with $i, j = 1, 2$.

3.2 Decays of electroweak SUSY particles

The second example concerns one-loop processes with external SUSY particles. A precise prediction of, e.g., SUSY production cross sections and decay branching ratios is necessary to obtain reliable bounds on the MSSM parameter space from LHC SUSY searches, to correctly interpret any possible signal at the LHC and to exploit the potential of a future e^+e^- collider such as the ILC, where measurements at the per-cent level will be possible.

Calculations of decay widths of SUSY particles in the cMSSM, using `MSSMCT.mod` (and its earlier versions) have been published in Refs. [17–19, 21–24]. Here we take one representative example from Ref. [18] that involves both a Higgs particle and a Dark Matter particle in the final state, the decay of the fourth neutralino to the lightest neutralino and the lightest Higgs boson, $\tilde{\chi}_4^0 \rightarrow \tilde{\chi}_1^0 h_1$. The parameters are given in Tab. 5. M_2 (the $SU(2)$ soft SUSY-breaking parameter) and μ (the Higgs mixing parameter) are chosen such that the values for $m_{\tilde{\chi}_1^\pm}$ and $m_{\tilde{\chi}_2^\pm}$ are fulfilled. Here the ambiguity in the hierarchy of M_2 and μ results in two scenarios: $\mu > M_2$ yields a higgsino-like $\tilde{\chi}_4^0$, denoted as S_h ; $\mu < M_2$ gives a gaugino-like $\tilde{\chi}_4^0$, denoted as S_g . $|M_1|$ (the absolute value of

the $U(1)$ soft SUSY-breaking parameter) is obtained from $|M_1| = \frac{5}{3} \tan^2 \theta_w M_2 \approx \frac{1}{2} M_2$, where φ_{M_1} is kept as a free parameter.

$\tan \beta$	M_{H^\pm}	$m_{\tilde{\chi}_2^\pm}$	$m_{\tilde{\chi}_1^\pm}$	$M_{\tilde{l}_L}$	$M_{\tilde{l}_R}$	A_l	$M_{\tilde{q}_L}$	$M_{\tilde{q}_R}$	A_q
20	160	600	350	300	310	400	1300	1100	2000

Table 5: MSSM parameters with all mass parameters are in GeV. M_{H^\pm} denotes the mass of the charged Higgs boson, $m_{\tilde{\chi}_{1,2}^\pm}$ are the chargino masses, $M_{\tilde{l}_L}$ and $M_{\tilde{l}_R}$ are the diagonal entries in the slepton mass matrices (taken to be universal for the three generations), A_l is the trilinear Higgs-slepton coupling; $M_{\tilde{q}_L}$, $M_{\tilde{q}_R}$ and A_q are the corresponding squark sector parameters (see Ref. [18] for details).

The absolute size of $\Gamma(\tilde{\chi}_4^0 \rightarrow \tilde{\chi}_1^0 h_1)$ in S_g and S_h is shown in the left plot of Fig. 2 as a function of φ_{M_1} ; separately shown are the tree-level results and the full one-loop calculation. A strong dependence of the absolute value of the decay width in both scenarios at the tree-level and at the one-loop level on this phase can be observed. The right plot shows the relative size of the one-loop corrections. Again a strong dependence of the size of those corrections on the phase of M_1 can be seen. Besides the default implementation of our renormalization scheme, denoted as S_g and S_h , we also show the results in an alternative on-shell scheme [18, 40] that differs in the treatment of the complex phases, denoted as \tilde{S}_g and \tilde{S}_h . It can be observed that the two schemes agree for real parameters ($\varphi_{M_1} = 0, \pi$) and show small differences (indicating the size of respective two-loop corrections) for complex parameters at or below the per-cent level.

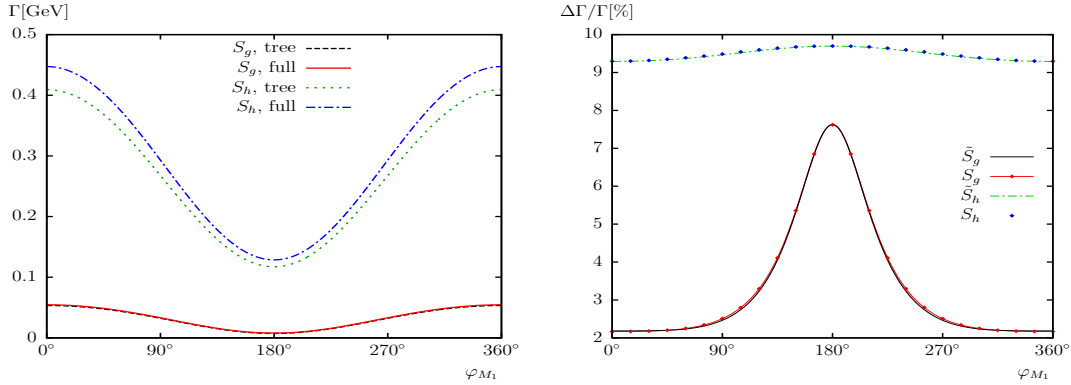


Figure 2: Left: absolute values of $\Gamma(\tilde{\chi}_4^0 \rightarrow \tilde{\chi}_1^0 h_1)$ in the scenarios S_g and S_h (see text) at the tree-level and the full one-loop level as a function of φ_{M_1} . Right: relative size of the one-loop corrections in our default renormalization scheme (S_g and S_h) and in an alternative scheme (\tilde{S}_g and \tilde{S}_h , see text) as a function of φ_{M_1} .

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